

THE WONDERFUL WORLD OF NEUTRINOS

Probir Roy
TIFR, Mumbai

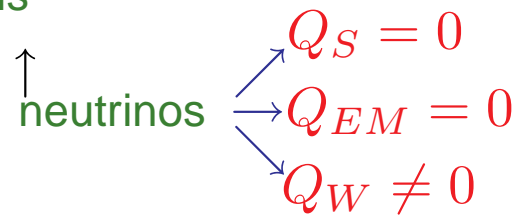
Kolkata Kolon Colloquium : SNBNCBS, December 1, 2004

“ . . . riddle wrapped in mystery inside an enigma ” Winston Churchill

- INTRODUCTION
- SOLAR NEUTRINOS
- ATMOSPHERIC NEUTRINOS
- REACTOR NEUTRINOS
- COSMOLOGICAL NEUTRINOS
- NEUTRINO FACTFILE
- LONG BASELINE NEUTRINO STUDIES
- CONCLUDING REMARKS

INTRODUCTION

elementary fermions: quarks and leptons



nucleus $\supset p, n$: mass ~ 1 GeV

$$c = 1 = \hbar$$

$$p = (uud), \quad n = (udd)$$

$$\begin{pmatrix} u \\ d \end{pmatrix} = \text{up quark with charge } +\frac{2}{3}|e|, \quad m_u \sim 3 \text{ MeV}$$
$$\begin{pmatrix} d \\ u \end{pmatrix} = \text{down quark with charge } -\frac{1}{3}|e|, \quad m_d \sim 7 \text{ MeV}$$

Two more doublets

$$\begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{matrix} m_c \sim 1.2 \text{ GeV} \\ m_s \sim 120 \text{ MeV} \end{matrix}$$

$$\begin{pmatrix} t \\ b \end{pmatrix} \quad \begin{matrix} m_t \sim 180 \text{ GeV} \\ m_b \sim 4.3 \text{ MeV} \end{matrix}$$

Three generations or flavours.

$$\begin{array}{l}
 \text{Also, } \begin{pmatrix} \nu_e \\ e \end{pmatrix} \quad m_e \sim 0.5 \text{ MeV} \\
 \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} \quad m_\mu \sim 105 \text{ MeV} \\
 m_\nu = ? \quad \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix} \quad m_\tau \sim 1.8 \text{ GeV}
 \end{array}$$

+ antiparticles \bar{f} for each f with the same mass.

Beta decay:

$$n \rightarrow p e^- \bar{\nu}_e$$

$$d(ud) \rightarrow u(ud) e^- \bar{\nu}_e$$

$$(\bar{u}d) : \pi^- \longrightarrow \nu^- \bar{\nu}_\mu \quad ; \quad \mu^- \longrightarrow e^- \bar{\nu}_e \nu_\mu$$

$$(u\bar{d}) : \pi^+ \longrightarrow \mu^+ \nu_\mu \quad ; \quad \mu^+ \longrightarrow e^+ \nu_e \bar{\nu}_\mu$$

$$Z \longrightarrow \tau^+ \tau^- \quad ; \quad \tau^+ \longrightarrow \mu^+ (e^+) \nu_\mu (\nu_e) \bar{\nu}_\tau$$

$$\tau^- \longrightarrow \mu^- (e^-) \bar{\nu}_\mu (\bar{\nu}_e) \nu_\tau$$

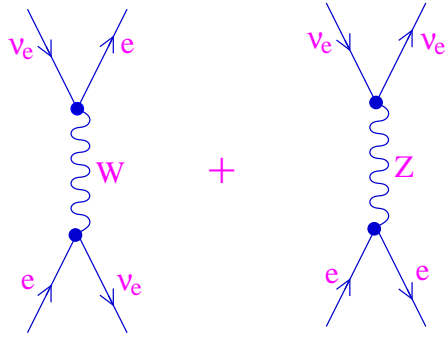
All sorts of neutrinos

Neutrino source	Description	Energy
Big Bang	Thermalized at 1.95°K . Undetected yet.	$\sim 10^{-4}$ eV
Stellar core	From nuclear reactions powering the star. Solar neutrinos detected.	~ 0.1 to 0 (10) MeV
Supernova	Detected from SN1987A .	10–30 MeV
Atmospheric	In a cosmic ray air-shower. Detected first in 1965 in Kolar Gold Fields, India.	sub-GeV to multi-GeV
Radioactivity on earth	Flux $\sim 10^6 \text{ cm}^{-2} \text{ sec}^{-1}$. Can be of geophysical use	\mathcal{O} (MeV)
Reactors	Antineutrinos from $n \rightarrow pe^- \bar{\nu}_e$ ICF in neutron-rich fissile nuclei.	~ 4 MeV
Accelerators	Neutrino beams at CERN, BNL, Fermilab, JHF.	sub-GeV to GeV

There can be additional cosmic sources of **UHE** neutrinos: **AGN, GRB.**

ICECUBE detector at South pole: 1 km^3 cubic lattice of phototubes.

Neutrinos may be Dirac particles, i.e. $\nu \neq \bar{\nu}$ or Majorana particles $\nu = \bar{\nu}$.

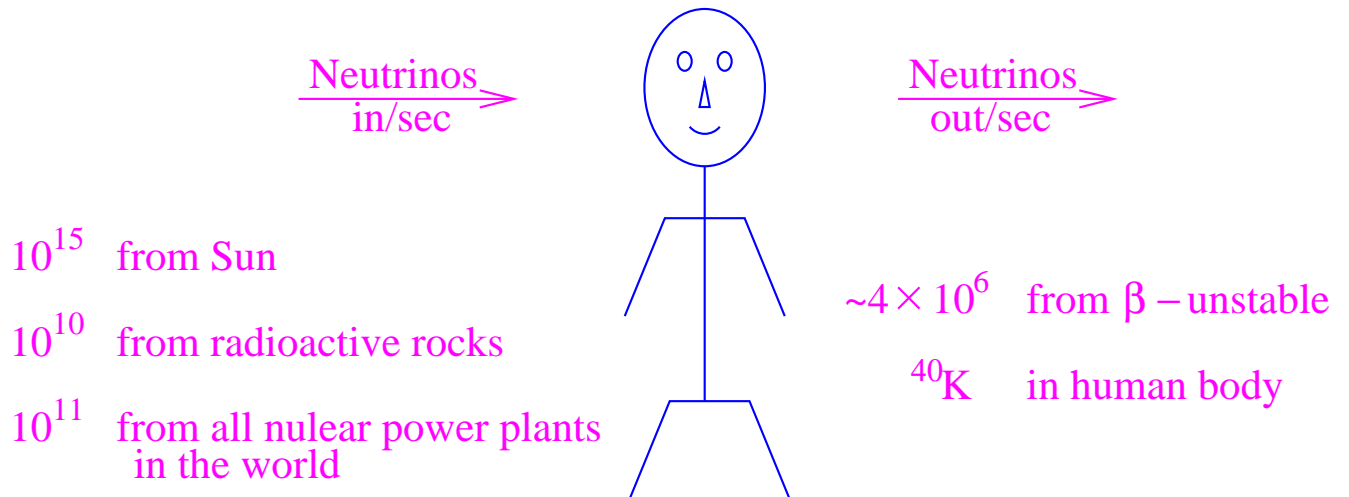


Neutrino scattering from matter electrons :

$\sigma \sim 10^{-44}$ cms for an MeV energy ν_e .

$$\lambda = 1/(N_e \sigma) \sim \pi / (2N_e G_F^2 m_e E_\nu)$$

Mfp of such ν in lead = 1 light year. Yet neutrino flavour conversion can be significantly enhanced in 1000 kms. of earth matter.



- In the Standard Model neutrinos are massless and only $\nu_L^{e,\mu,\tau} = \frac{1}{2}(1 - \gamma_5)\nu^{e,\mu,\tau}$ exist
- Almost all extensions of the Standard Model predict nonzero neutrino masses

Dirac ν has four components $\nu_L, \nu_L^C, N_R, N_R^C$.

$$\nu = \nu_L + N_R.$$

Dirac mass term $m_D \bar{\nu}\nu = m_D(\bar{\nu}_L N_R + \bar{N}_R \nu_L)$

Question: Why is $m_\nu \ll m_f \quad \forall f$.

See-saw Mechanism

$$\begin{array}{c} \nu_L \\ \bar{N}_R \end{array} \begin{array}{cc} \nu_L & \bar{N}_R \\ \left[\begin{array}{cc} 0 & m_D \\ m_D & M_N \end{array} \right] \end{array} \rightarrow \begin{pmatrix} \sim \frac{m_D^2}{M_N} & 0 \\ 0 & \sim M_N \end{pmatrix}$$

$M_N \gg m_D = \mathcal{O}(m_f) \quad \uparrow$
 large Majorana
 mass for N_R

Now $m_\nu \sim m_D^2/M_N \ll m_D$.

Other proposed mechanisms for the smallness of $m_\nu \sim \frac{m_f}{(V_\delta)^{1/2}}$
 in terms of extra compactified dimensions.

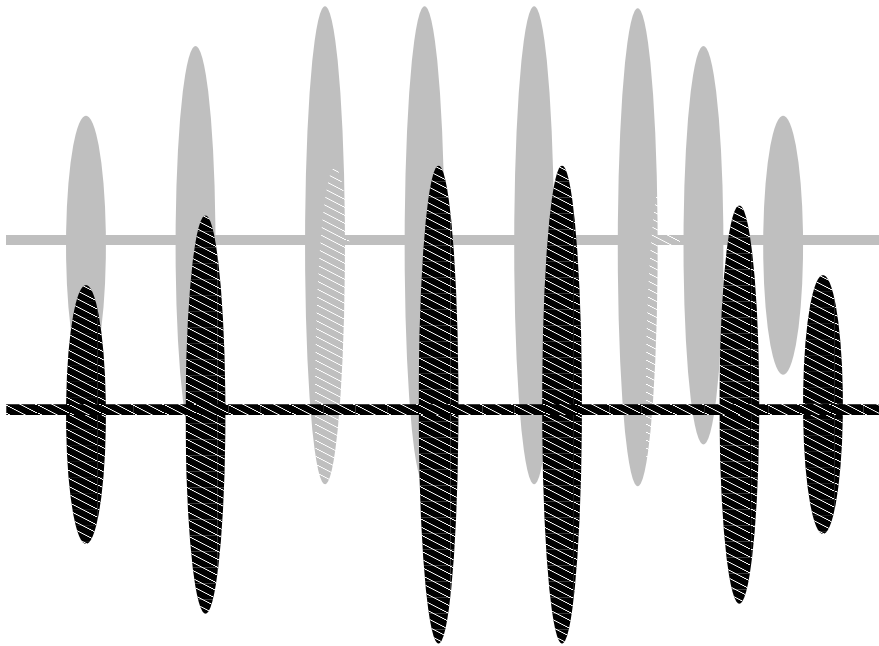
Neutrino Oscillations

ν weak eigenstates may be different from ν mass eigenstates.

2 flavours situation, mixing angle θ :

$$\nu_e = \nu_1 \cos \theta + \nu_2 \sin \theta$$

$$\nu_f = -\nu_1 \sin \theta + \nu_2 \cos \theta$$



Weak decay produces distinct weak eigenstate $\nu(0) = \nu_e$, say

$$\nu(t) = \nu_1 e^{-iE_1 t} \cos \theta + \nu_2 e^{-iE_2 t} \sin \theta$$

Ultrarelativistically, $p \sim E - \frac{m^2}{2E}$.

$$P[\nu_e \rightarrow \nu_f; L] = \sin^2 2\theta \sin^2 \frac{E_2 - E_1}{2} t.$$

$$P[\nu_e \rightarrow \nu_f; L] \simeq \sin^2 2\theta \sin^2 \frac{\delta m_{21}^2 L}{4E}$$

$$\simeq \sin^2 2\theta \sin^2 \frac{1.27(\delta m_{21}/eV)^2 (L/km)}{E/GeV}$$

SOLAR NEUTRINOS

ν_e 's with sub-MeV to a few MeV energies from the solar core

Main nuclear reactions

<i>Reaction</i>	<i>Max. flux in $m^{-2}s^{-1}$ from SSM</i>	<i>E_ν</i>
$pp \rightarrow de^+\nu_e$	$\sim 6 \times 10^{14}$	<i>0 to 0.420 MeV</i>
$pep \rightarrow d\nu_e$	$\sim 1.5 \times 10^{12}$	<i>1.8 MeV</i>
${}^7Be e \rightarrow {}^7Li \nu_e$	$\sim 4.75 \times 10^{13}$	<i>0.38 MeV & 0.87 MeV</i>
${}^8B \rightarrow {}^8B^* e^+ \nu_e$	$\sim 5 \times 10^{10}$	<i>0 to 14.6 MeV</i>
${}^{13}N \rightarrow {}^{12}C e^+ \nu_e$	$\sim 5.5 \times 10^{12}$	<i>0 to 1.25 MeV</i>
${}^{15}O \rightarrow {}^{15}N e^+ \nu_e$	$\sim 4.8 \times 10^{12}$	<i>0 to 1.75 MeV</i>
${}^{17}F \rightarrow {}^{17}O e^+ \nu_e$	$\sim 5.5 \times 10^{10}$	<i>0 to 1.75 MeV</i>

Detection (deep underground)

(1) Radiochemical



DAVIS

Nobel 2002



*Atoms counted
using Auger effect.*



(2) *Realtime scattering in water* $\nu_e \rightarrow \nu_e e$.

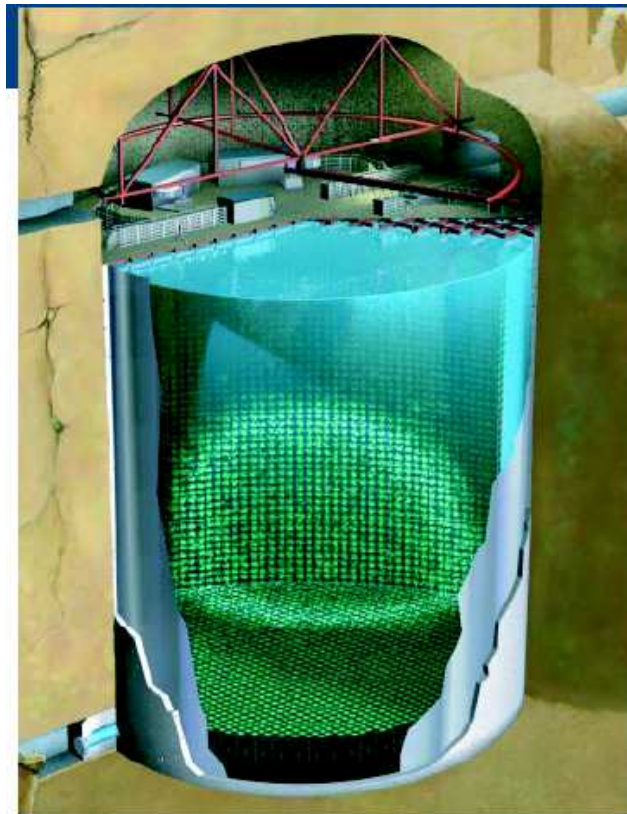
KOSHIBA

Works less efficiently

Nobel 2002

for ν_μ and ν_τ : $\nu_{\mu,\tau} e \rightarrow \nu_{\mu,\tau} e$

Directionality (to Sun) via Cerenkov cones.

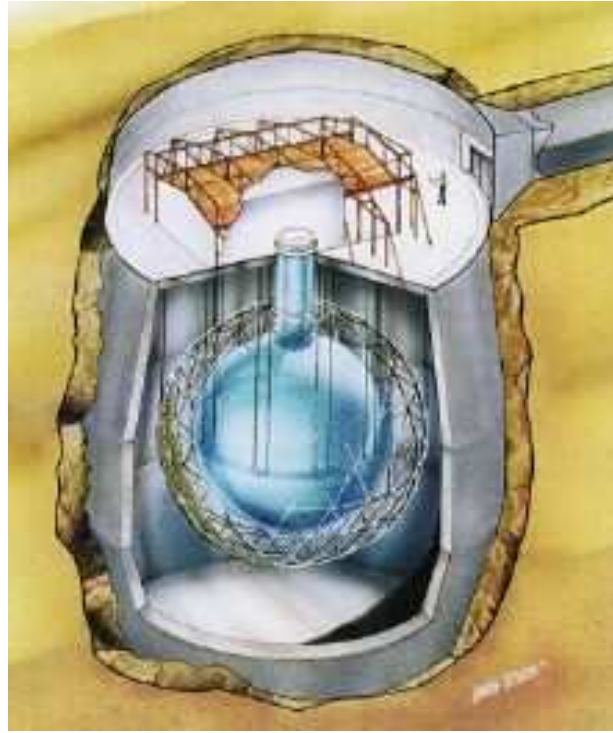


(3) *Realtime scattering in heavy water* D_2O (SNO)

$$\nu_e d \rightarrow e p p$$

$$\nu_x d \rightarrow \nu_x p n$$

$$x = e, \mu, \tau$$



Reduced ν_e fluxes detected as compared to SSM expectations.

Explanation: flavour conversion $\nu_e \rightarrow \nu_{\mu,\tau}$.

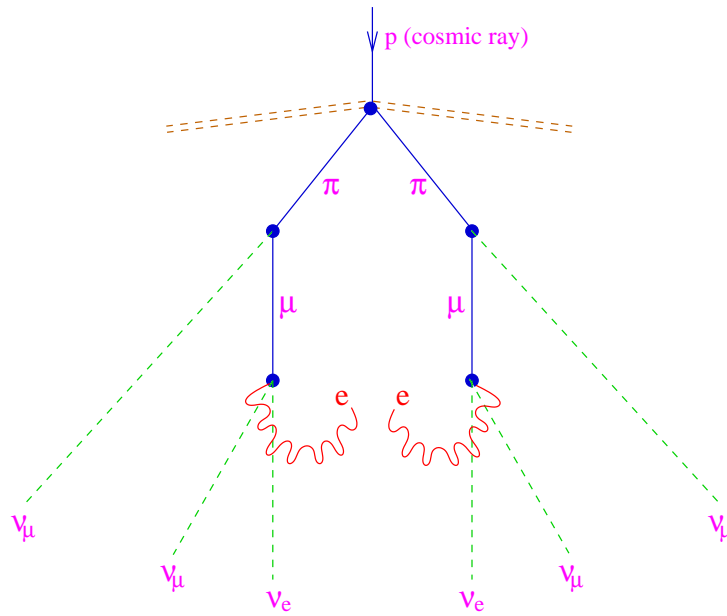
$$\text{SNO} = \frac{\text{observed } \nu_e \text{ flux}}{\text{observed } \nu_x \text{ flux}} = \frac{(\nu_e - \text{flux})_{\text{Earth}}}{(\nu_e - \text{flux})_{\text{Sun}}}$$

$$P_{ee} = 1 - \sin^2 2\theta_{\odot} \sin^2 \frac{1.27(\delta m_{21}^2 / eV^2)(L/km)}{E/GeV}$$

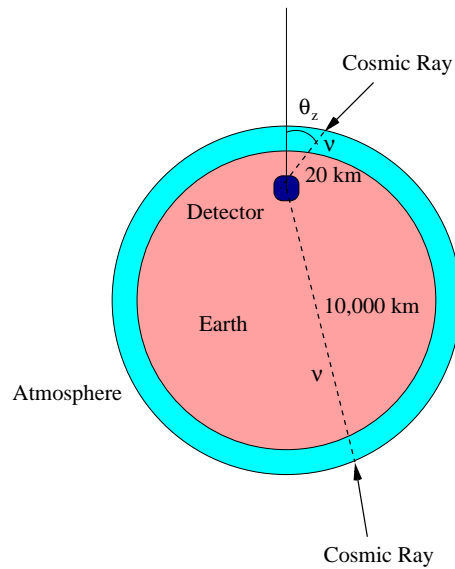
+ Resonant enhancement of flavour conversion in solar medium MSW

$$\implies \delta m_{21}^2 = (6.8 \pm 0.8) \times 10^{-5} eV^2, \theta_S = 32.5^\circ \pm 2.5^\circ.$$

ATMOSPHERIC NEUTRINOS

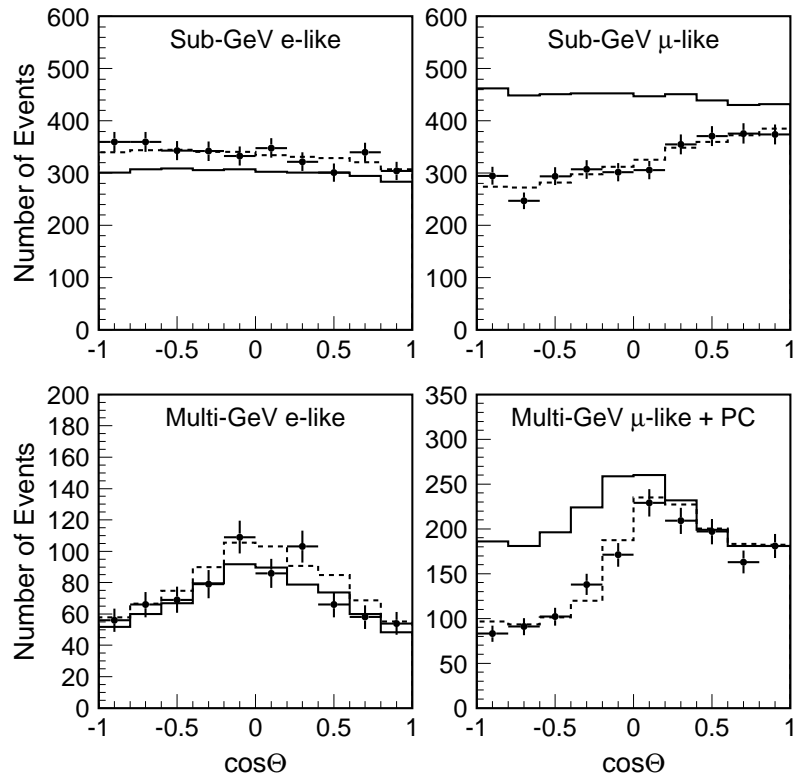


Roughly, expected $\frac{N(\nu_{\mu})}{N(\nu_e)} \simeq 2$. Seen: almost equal nos.



sub-GeV: $\frac{(N_{\mu}/N_e)_{data}}{(N_{\mu}/N_e)_{expected}} = 0.652 \pm 0.019 \pm 0.051$

multi-GeV μ -like: $\frac{N_{\mu}^{up} - N_{\mu}^{down}}{N_{\mu}^{up} + N_{\mu}^{down}} = -0.296 \pm 0.032 \pm 0.01$



Best interpretation: $\nu_\mu \leftrightarrow \nu_\tau$ oscillation

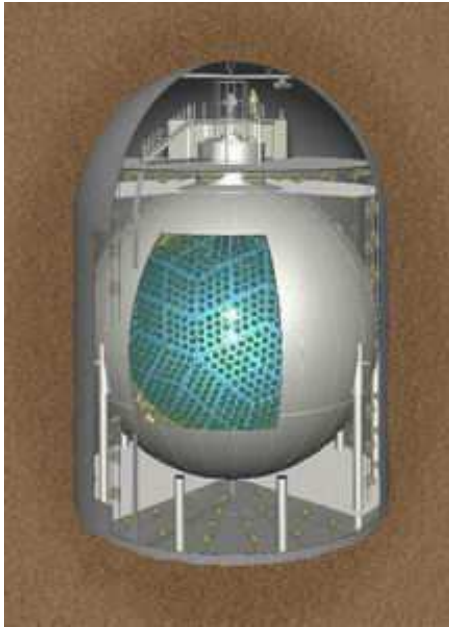
$$|\delta m_{32}^2| = (2.6 \pm 0.4) \times 10^{-3} eV^2$$

$$\theta_A = 45^\circ \pm 5^\circ.$$

Indian initiative: India-based Neutrino Observatory (INO)

Underground calorimetric detector: magnetized iron layers with RPC's.

REACTOR NEUTRINOS



KamLAND expt:

$\bar{\nu}_e$ beams from 50 reactors between 150 and 350 kms into 1 kton liquid scintillator detector at Kamioka. Detection via $\bar{\nu}_e + p \rightarrow e^+ + n$. Measured flux vs. input calculated flux clearly supports δm_{21}^2 and θ_{\odot} values.

CHOOZ expt:

$\bar{\nu}_e$'s from two cores in a French nuclear power plant.

No depletion found at $L \sim 1\text{km} \implies |\theta_{13}| < 13^\circ$.

Double-CHOOZ expt. being undertaken with one more "near" detector at $L \sim 100 - 200\text{m}$ from the cores.

Hopes to improve limit on θ_{13} by factor 3.

New KASKA (Kashiwazaki-Kariwa complex) expt in Japan, Krasnoyarsk in Russia and Daya Bay in China.

COSMOLOGICAL NEUTRINOS

Generated from the Big Bang

$$N(p) = \frac{1}{e^{\frac{p}{kT_c}} + 1}$$

Decoupled about 1 min after BB. HDM today

About 330 ($\nu, \bar{\nu}$)'s/cm³.

$$\Omega = \rho / \rho_c$$

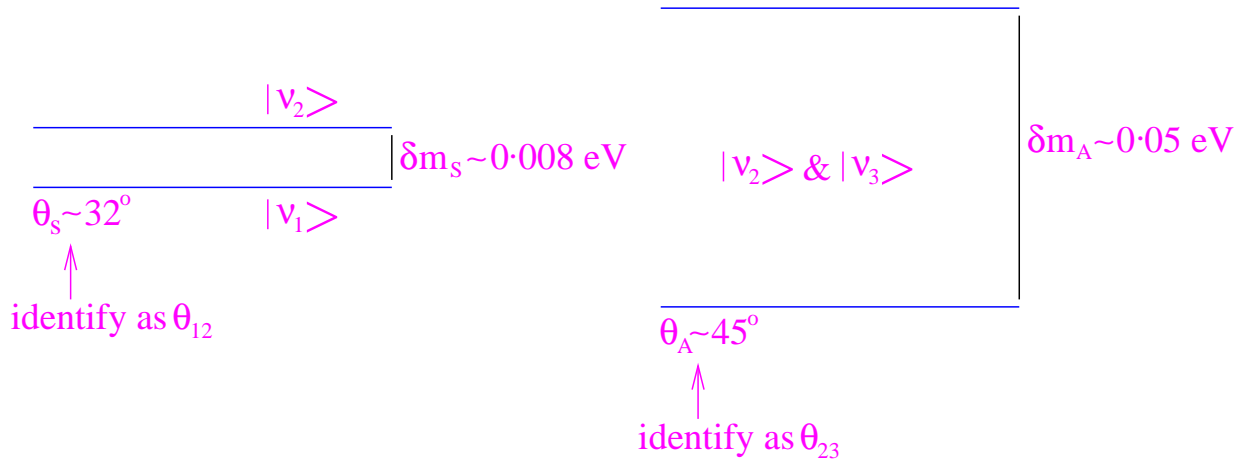
$$\Omega_{\nu, \bar{\nu}} h^2 = \frac{\sum_{\nu} m_{\nu}}{91.5 \text{ eV}},$$

$$H_0 = 100 h \text{ km sec}^{-1} \text{ mpc}^{-1} \text{ and } h \sim 0.7$$

WMAP probe into different components of Ω has bounded Ω_{HDM} from above

$$\Rightarrow \sum_{\nu} m_{\nu} < 0.71 \text{ eV.}$$

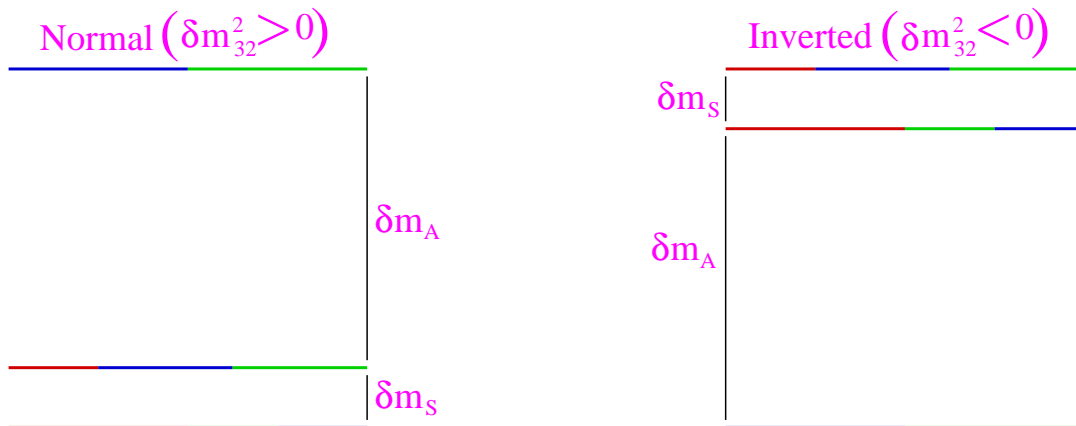
NEUTRINO FACTFILE



$$|\theta_{13}| < 13^\circ$$

Two possible types of ordering

$$\delta m_{ij}^2 = m_i^2 - m_j^2.$$



LONG BASELINE NEUTRINO STUDIES

- *K2K*
- *T2K*
- *MINOS*
- *CNGS*
 - *ICARUS*
 - *OPERA*
- *MINERVA*
- *NOVA*

Accelerator neutrino ($\nu_\mu, \bar{\nu}_\mu$ or β -beams $\nu_e, \bar{\nu}_e$) beam + near detector + far detector

<i>Project</i>	<i>Accelerator</i>	<i>Location of far detector</i>	<i>Dist. kms.</i>	$\langle E_\nu \rangle$ <i>GeV</i>	<i>Status</i>
<i>K2K</i>	<i>KEK</i> <i>p</i> <i>synchrotron</i>	<i>Kamioka mines</i>	<i>250</i>	<i>1.4</i>	<i>Started April '99</i>
<i>Minos</i>	<i>Fermilab</i> <i>Main injector</i>	<i>Soudan mine</i>	<i>730</i>	<i>3 ± 1</i>	<i>To start in 2005</i>
<i>CNGS</i>	<i>CERN</i> <i>450 GeV SPS</i>	<i>Gransasso Lab</i>	<i>732</i>	<i>a few GeV</i>	<i>– do –</i>
<i>T2K</i>	<i>p</i> <i>accelerator</i> <i>JHF, Tokai</i>	<i>Kamloka mines</i>	<i>295</i>	<i>≤ 1 GeV</i>	<i>?</i>
<i>Minerva</i> <i>Nova</i>	<i>Fermilab main injector</i> <i>Superberm ?</i>	<i>Soudan 15 km off-axis surface</i>	<i>732+</i>	<i>$1 \simeq 2$ GeV</i>	<i>?</i>

superbeams
neutrino factories

Our result

*B. Brahmachari, S. Choubey & PR: Nucl. Phys. **B671** (2003) 483.*
*S. Choubey & PR: Phys. Rev. Lett. **93** (2004) 021803.*

$|\nu_\alpha\rangle =$ neutrino flavour states ($\alpha = e, \mu, \tau$) participating in EW gauge interactions.

$|\nu_i\rangle =$ neutrino mass eigenstates with masses $m_i (i, 1, 2, 3)$

$$|\nu_\alpha\rangle = U_{\alpha i} |\nu_i\rangle$$



unitary transformation

$$U \equiv U_{PMNS} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix}$$

$$\left. \begin{array}{l} |\nu_1\rangle \rightarrow |\nu_e\rangle \\ |\nu_2\rangle \rightarrow |\nu_\mu\rangle \\ |\nu_3\rangle \rightarrow |\nu_\tau\rangle \end{array} \right\} \text{when } U \rightarrow I : \text{ PMNS convention}$$

Propagation in vacuum

$$|\nu_i\rangle_L = e^{ip_i L} |\nu_i\rangle_0$$

$$|\nu_\alpha\rangle_L = U_{\alpha i} |\nu_i\rangle_L = U_{\alpha i} U_{\beta i}^* e^{ip_i L} |\nu_\beta\rangle_0$$

$$S_{\alpha \rightarrow \beta}(L) = \langle \nu_\beta | \nu_\alpha \rangle_0 = U_{\alpha i} U_{\beta i}^* e^{ip_i L}.$$

$$P[\nu_\alpha(0) \rightarrow \nu_\beta(L)] = |S_{\alpha \rightarrow \beta}(L)|^2.$$

Ultrarelativistically, $p \sim E - \frac{m^2}{2E}$. Define $\Delta_{ij} \equiv \frac{m_i^2 - m_j^2}{4E} L$.

In the approximation $|\delta m_{32}|^2 \gg |\delta m_{21}^2|$ one can show that

$$\begin{aligned} P[\nu_\mu(0) \rightarrow \nu_\mu(L)]_{vac} &\simeq 1 - 4|U_{\mu 3}|^2(1 - |U_{\mu 3}|^2) \sin^2 \Delta_{31} \\ &= P[\bar{\nu}_\mu(0) \rightarrow \bar{\nu}_\mu(L)]_{vac} \end{aligned}$$

ν_μ survival probability in vacuum, minimum for $|U_{\mu 3}| = 1/\sqrt{2}$.

ν_μ flavour conversion probability maximal for $|U_{\mu 3}| = 1/\sqrt{2}$.

$|U_{\mu 3}|$ known to be $\sim 1/\sqrt{2}$.

$|\frac{1}{\sqrt{2}} - |U_{\mu 3}||$ is the deviation from maximality difficult to measure at a minimum.

CONCLUDING REMARKS

A glimpse of the world of neutrinos.

Only indication of new physics beyond the Standard Model.

Not covered

- *Geothermal neutrinos*
- *Neutrino astronomy*
- *CP violation in neutrino sector*